## Convergent FE scheme for the two-fluid MHD equation

Markus Klein

Institut of Mathematics
University of Tuebingen

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## Outline

(1) Introduction, prelimitaries
(2) Discretization, Convergence

- Continuous Galerkin Approach
- Discontinuous Galerkin Approach
(3) Some words on implementation

4 Summary

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## Application

Production of Al from $\mathrm{Al}_{2} \mathrm{O}_{3}$ by Electrolysis (cf. [Gerbeau et al., 2006, chap 6]):

(chemie.uni-freiburg.de)

## Existence of weak solutions, [Gerbeau et al., 2006]

Under certain assumptions on the initial data, there exists a weak solution $\boldsymbol{u} \in L^{\infty}\left(0, T ; \boldsymbol{L}^{2} \cap\{\operatorname{div} \boldsymbol{u}=0\right.$ weakly $\left.\}\right) \cap L^{2}\left(0, T ; \boldsymbol{W}_{0}^{1,2} \cap\{\operatorname{div} \boldsymbol{u}=0\right.$ a.e. $\left.\}\right)$, $\boldsymbol{b} \in L^{\infty}\left(0, T ; \boldsymbol{L}^{2} \cap\{\operatorname{div} \boldsymbol{b}=0\right.$ weakly $\left.\}\right) \cap L^{2}(0, T ; \boldsymbol{X})$, $\rho \in L^{\infty}((0, T) \times \Omega) \cap \mathcal{C}\left([0, T], L^{p}\right)$ which holds the property

$$
\mid\{x \in \Omega: \alpha \leq \rho(x, t) \leq \beta \mid \text { is constant in time for all } 0 \leq \alpha \leq \beta<\infty .
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## Remark

- Since we know $\rho_{0}$, the fluids are moving within $\Omega$ !
- If $\rho^{n}$ solves discretized eq. and " $\rho$ " $\rightarrow \rho$ ", the property holds for $\rho^{n}$ appr.


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## Incredients for the proof

- Typical technical arguments related to one-fluid MHD.
- Use of DiPerna-Lions compactness, cf. [DiPerna and Lions, 1989] (for passing to a limit of $\rho$ ).


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## General setup

- $\mathcal{T}_{h}$ quasi-uniform trangulation of polyhedral domain $\Omega \subseteq \mathbb{R}^{d}(d=2,3)$.
- $V_{h}:=\left\{\xi \in \mathcal{C}(\bar{\Omega}): \xi \in P_{1}(T) \forall T \in \mathcal{T}_{h}\right\}$ FE space w.r.t. $\rho$.
- $\boldsymbol{V}_{h} \subseteq \boldsymbol{W}_{0}^{1,2}$ FE space w.r.t. $\boldsymbol{u}$ and $L_{h} \subseteq L_{0}^{2}$ FE space w.r.t. $P$, s.t. $\left(\boldsymbol{V}_{h}, L_{h}\right)$ holds inf-sup cond.
- $\boldsymbol{C}_{h}:=\left\{\boldsymbol{\psi} \in \boldsymbol{H}\right.$ (curl) : $\psi \in \boldsymbol{\mathcal { N }}_{j}$ for some $\left.j \geq 1\right\}$ (Nedelec) FE space w.r.t $\boldsymbol{b}$.
- $S_{h} \subseteq W^{1,2} \cap L_{0}^{2}$ s.t. inf-sup cond. holds for $\left(\boldsymbol{C}_{h}, S_{h}\right)$.


## Idea of discretization

Reformulate:

$$
(\rho \boldsymbol{u})_{t}+\operatorname{div}(\rho \boldsymbol{u} \otimes \boldsymbol{u})=\frac{1}{2}\left(\rho \boldsymbol{u}_{t}+(\rho \boldsymbol{u} \cdot \nabla) \boldsymbol{u}+(\rho \boldsymbol{u})_{t}+\operatorname{div}(\rho \boldsymbol{u} \otimes \boldsymbol{u})\right)
$$

(true, since $\left.\rho_{t}+\operatorname{div}(\rho \mathbf{u})=0\right)$.

## Algorithm ((Scheme A), [Baňas and Prohl, 2010])

Find $\left(\boldsymbol{U}^{n}, P^{n}, \boldsymbol{B}^{n}, R^{n}, \rho^{n}\right)$ s.t. for all $(\chi, \boldsymbol{W}, \boldsymbol{\psi})$ :

$$
\begin{aligned}
& 0=\left(d_{t} \rho^{n}, \chi\right)_{n}+\left(\boldsymbol{U}^{n} \cdot \nabla \rho^{n}, \chi\right)+\frac{1}{2}\left(\operatorname{div}\left(\boldsymbol{U}^{n}\right) \rho^{n}, \chi\right) \\
&\left(\rho^{n-1} \boldsymbol{g}^{n}, \boldsymbol{W}\right)= \frac{1}{2}\left\{\left(\rho_{+}^{n-1} d_{t} \boldsymbol{U}^{n}, \boldsymbol{W}\right)_{*}+\left(d_{t}\left(\rho_{+}^{n} \boldsymbol{U}^{n}\right), \boldsymbol{W}\right)_{*}+\left(\left(\rho^{n-1} \boldsymbol{U}^{n-1} \cdot \nabla\right) \boldsymbol{U}^{n}, \boldsymbol{W}\right)-\left(\left(\rho^{n-1} \boldsymbol{U}^{n-1} \cdot \nabla\right) \boldsymbol{W}, \boldsymbol{U}^{n}\right)\right\} \\
&+\left(\eta^{n-1} \boldsymbol{D}\left(\boldsymbol{U}^{n}\right), \boldsymbol{D}(\boldsymbol{W})\right)-\left(P^{n}, \operatorname{div} \boldsymbol{W}\right)+\frac{1}{\bar{\mu}}\left(\boldsymbol{B}^{n-1} \times \operatorname{curl} \boldsymbol{B}^{n}, \boldsymbol{W}\right) \\
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\left(\rho^{n-1} \boldsymbol{g}^{n}, \boldsymbol{W}\right)= & \frac{1}{2}\left\{\left(\rho_{+}^{n-1} d_{t} \boldsymbol{U}^{n}, \boldsymbol{W}\right)_{*}+\left(d_{t}\left(\rho_{+}^{n} \boldsymbol{U}^{n}\right), \boldsymbol{W}\right)_{*}+\left(\left(\rho^{n-1} \boldsymbol{U}^{n-1} . \nabla\right) \boldsymbol{U}^{n}, \boldsymbol{W}\right)-\left(\left(\rho^{n-1} \boldsymbol{U}^{n-1} . \nabla\right) \boldsymbol{W}, \boldsymbol{U}^{n}\right)\right\} \\
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for some non negative constants $\alpha$
$\left(\eta^{n-1}:=\eta\left(\rho^{n-1}\right)\right.$ and $\left.\xi^{n-1}:=\xi\left(\rho^{n-1}\right)\right)$

- $\alpha$ term $\Rightarrow$ M-Matrix property of the Scheme


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- $\alpha$ term $\Rightarrow$ M-Matrix property of the Scheme
- $\beta_{1}$ term $\Rightarrow \boldsymbol{L}^{2}$-strong convergence of $\operatorname{div} \boldsymbol{U}^{n}$


## Algorithm ((Scheme A), [Bañas and Prohl, 2010])

Find $\left(\boldsymbol{U}^{n}, P^{n}, \boldsymbol{B}^{n}, R^{n}, \rho^{n}\right)$ s.t. for all $(\chi, \boldsymbol{W}, \boldsymbol{\psi})$ :

$$
\begin{aligned}
0= & \left(d_{t} \rho^{n}, \chi\right)_{h}+\left(\boldsymbol{U}^{n} \cdot \nabla \rho^{n}, \chi\right)+\frac{1}{2}\left(\operatorname{div}\left(\boldsymbol{U}^{n}\right) \rho^{n}, \chi\right)+\alpha h^{\alpha}\left(\nabla \rho^{n}, \nabla \chi\right) \\
\left(\rho^{n-1} \boldsymbol{g}^{n}, \boldsymbol{W}\right)= & \frac{1}{2}\left\{\left(\rho_{+}^{n-1} d_{t} \boldsymbol{U}^{n}, \boldsymbol{W}\right)_{*}+\left(d_{t}\left(\rho_{+}^{n} \boldsymbol{U}^{n}\right), \boldsymbol{W}\right)_{*}+\left(\left(\rho^{n-1} \boldsymbol{U}^{n-1} . \nabla\right) \boldsymbol{U}^{n}, \boldsymbol{W}\right)-\left(\left(\rho^{n-1} \boldsymbol{U}^{n-1} . \nabla\right) \boldsymbol{W}, \boldsymbol{U}^{n}\right)\right\} \\
& +\left(\eta^{n-1} \boldsymbol{D}\left(\boldsymbol{U}^{n}\right), \boldsymbol{D}(\boldsymbol{W})\right)-\left(P^{n}, \operatorname{div} \boldsymbol{W}\right)+\frac{1}{\bar{\mu}}\left(\boldsymbol{B}^{n-1} \times \operatorname{curl} \boldsymbol{B}^{n}, \boldsymbol{W}\right) \\
& +\beta_{2} h^{-\beta_{1}}\left(\operatorname{div} \boldsymbol{U}^{n}, \operatorname{div} \boldsymbol{W}\right)+\beta_{2} h^{\beta_{2}}\left(\nabla d_{t} \boldsymbol{U}^{n}, \nabla \boldsymbol{W}\right) \\
0= & \left(d_{t} \boldsymbol{B}^{n}, \boldsymbol{\psi}\right)+\frac{1}{\bar{\mu}}\left(\frac{1}{\xi^{n-1}} \operatorname{curl} \boldsymbol{B}^{n}, \operatorname{curl} \psi\right)-\left(\boldsymbol{U}^{n} \times \boldsymbol{B}^{n-1}, \operatorname{curl} \psi\right)-\left(\nabla R^{n}, \boldsymbol{\psi}\right)
\end{aligned}
$$

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$$
\begin{aligned}
0= & \left(d_{t} \rho^{n}, \chi\right)_{h}+\left(\boldsymbol{U}^{n} . \nabla \rho^{n}, \chi\right)+\frac{1}{2}\left(\operatorname{div}\left(\boldsymbol{U}^{n}\right) \rho^{n}, \chi\right)+\alpha h^{\alpha}\left(\nabla \rho^{n}, \nabla \chi\right) \\
\left(\rho^{n-1} \boldsymbol{g}^{n}, \boldsymbol{W}\right)= & \frac{1}{2}\left\{\left(\rho_{+}^{n-1} d_{t} \boldsymbol{U}^{n}, \boldsymbol{W}\right)_{*}+\left(d_{t}\left(\rho_{+}^{n} \boldsymbol{U}^{n}\right), \boldsymbol{W}\right)_{*}+\left(\left(\rho^{n-1} \boldsymbol{U}^{n-1} \cdot \nabla\right) \boldsymbol{U}^{n}, \boldsymbol{W}\right)-\left(\left(\rho^{n-1} \boldsymbol{U}^{n-1} . \nabla\right) \boldsymbol{W}, \boldsymbol{U}^{n}\right)\right\} \\
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0= & \left(d_{t} \boldsymbol{B}^{n}, \boldsymbol{\psi}\right)+\frac{1}{\bar{\mu}}\left(\frac{1}{\xi^{n-1}} \operatorname{curl} \boldsymbol{B}^{n}, \operatorname{curl} \psi\right)-\left(\boldsymbol{U}^{n} \times \boldsymbol{B}^{n-1}, \operatorname{curl} \psi\right)-\left(\nabla R^{n}, \boldsymbol{\psi}\right)
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- $\beta_{3}$ term $\Rightarrow$ strong $\boldsymbol{L}^{2}$-convergence of $\boldsymbol{U}^{n}$ (for $d=3$; for $d=2$ this is free due Sobolev imbeddings)


## Lemma (Existence, energy law, maximum principle, [Baňas and Prohl, 2010])

Let $\alpha, \beta_{1}, \beta_{2}, \beta_{3} \geq 0, \sqrt{\beta_{2}} h^{\beta_{2}} 2\left\|\nabla \boldsymbol{U}^{0}\right\| \leq C$. Then there exists a solution $\left(\boldsymbol{U}^{n}, \boldsymbol{B}^{n}, \rho^{n}, P^{n}, R^{n}\right)$ of the numerical scheme, which holds the discrete energy law

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\begin{aligned}
\int_{\Omega} \rho^{n-1} \boldsymbol{g}^{n} \boldsymbol{U}^{n}= & \frac{1}{2} \partial_{t}\left(\left\|\sqrt{\rho_{+}^{n}} \boldsymbol{U}^{n}\right\|_{*}^{2}+\beta_{2} h^{\beta_{2}}\left\|\nabla \boldsymbol{U}^{n}\right\|^{2}+\frac{1}{\bar{\mu}}\left\|\boldsymbol{B}^{n}\right\|^{2}\right)+\left\|\sqrt{\eta^{n-1}} \boldsymbol{D}\left(\boldsymbol{U}^{n}\right)\right\|^{2}+\beta_{1} h^{-\beta_{1}}\left\|\operatorname{div} \boldsymbol{U}^{n}\right\|^{2} \\
& +\beta_{3} h^{\beta_{3}}\left\|\Delta_{h} \boldsymbol{U}^{n}\right\|^{2}+\frac{1}{\bar{\mu}^{2}}\left\|\frac{1}{\sqrt{\xi^{n-1}}} \operatorname{curl} \boldsymbol{B}^{n}\right\|^{2}+\frac{k}{2}\left(\left\|\sqrt{\rho_{+}^{n-1}} d_{t} \boldsymbol{U}^{n}\right\|_{*}^{2}+\beta_{2} h^{\beta_{2}}\left\|\nabla d_{t} \boldsymbol{U}^{n}\right\|^{2}+\left\|d_{t} \boldsymbol{B}^{n}\right\|^{2}\right) \\
0= & \frac{1}{2} d_{t}\left\|\rho^{n}\right\|_{h}^{2}+\frac{k}{2}\left\|d_{t} \rho^{n}\right\|_{h}^{2}+\alpha h^{\alpha}\left\|\nabla \rho^{n}\right\|^{2}
\end{aligned}
$$

Let $V_{h} \cap L_{0}^{2} \subseteq L_{h}, \mathcal{T}_{h}$ be a strongly acute triangulation, $\alpha, \beta_{>} 0$ and $0<\alpha+\frac{\beta_{2}}{2}<\frac{6-d}{6}$. Then $0<\bar{\rho}_{1} \leq \rho^{n} \leq \bar{\varrho}_{2}<\infty$ (discrete maximum principle).

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& +\beta_{3} h^{\beta_{3}}\left\|\Delta_{h} \boldsymbol{U}^{n}\right\|^{2}+\frac{1}{\bar{\mu}^{2}}\left\|\frac{1}{\sqrt{\xi^{n-1}}} \operatorname{curl} \boldsymbol{B}^{n}\right\|^{2}+\frac{k}{2}\left(\left\|\sqrt{\rho_{+}^{n-1}} d_{t} \boldsymbol{U}^{n}\right\|_{*}^{2}+\beta_{2} h^{\beta_{2}}\left\|\nabla d_{t} \boldsymbol{U}^{n}\right\|^{2}+\left\|d_{t} \boldsymbol{B}^{n}\right\|^{2}\right), \\
0= & \frac{1}{2} d_{t}\left\|\rho^{n}\right\|_{h}^{2}+\frac{k}{2}\left\|d_{t} \rho^{n}\right\|_{h}^{2}+\alpha h^{\alpha}\left\|\nabla \rho^{n}\right\|^{2} .
\end{aligned}
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$0<\bar{\rho}_{1} \leq \rho^{n} \leq \bar{\varrho}_{2}<\infty$ (discrete maximum principle).

## Sketch of the proof

- Suppose

$$
\max _{0 \leq \ell \leq n-1}\left(\left\|\rho^{\ell}\right\|_{h}^{2}+\left\|\sqrt{\rho_{+}^{\ell}} \boldsymbol{U}^{\ell}\right\|_{*}^{2}+\beta_{2} h^{\beta_{2}}\left\|\nabla \boldsymbol{U}^{\ell}\right\|^{2}+\frac{1}{\bar{\mu}}\left\|\boldsymbol{B}^{\ell}\right\|^{2}\right) \leq C,
$$

(fullfilled for $n=1$ by asumption).

- Define $\mathcal{F}^{n}([\rho, \boldsymbol{U}, \boldsymbol{B}],[\chi, \boldsymbol{W}, \boldsymbol{\psi}]):=$ Scheme $A$ - right hand side (terms with $R^{n}, P^{n}$ vanish).
- Show: $\mathcal{F}^{n}([\rho, \boldsymbol{U}, \boldsymbol{B}],[\rho, \boldsymbol{U}, \boldsymbol{B}]] \geq 0 \Rightarrow \exists \rho^{n}, \boldsymbol{U}^{n}, \boldsymbol{B}^{n}$ with Brouwer. Boundness of $\rho^{\ell}$, etc. is fullfilled for $n+1$ by Brouwer.


## Lemma (Existence, energy law, maximum principle, [Baňas and Prohl, 2010])

Let $\alpha, \beta_{1}, \beta_{2}, \beta_{3} \geq 0, \sqrt{\beta_{2}} h^{\beta_{2}} 2\left\|\nabla \boldsymbol{U}^{0}\right\| \leq C$. Then there exists a solution $\left(\boldsymbol{U}^{n}, \boldsymbol{B}^{n}, \rho^{n}, P^{n}, R^{n}\right)$ of the numerical scheme, which holds the discrete energy law

$$
\begin{aligned}
\int_{\Omega} \rho^{n-1} \boldsymbol{g}^{n} \boldsymbol{U}^{n}= & \frac{1}{2} \partial_{t}\left(\left\|\sqrt{\rho_{+}^{n}} \boldsymbol{U}^{n}\right\|_{*}^{2}+\beta_{2} h^{\beta_{2}}\left\|\nabla \boldsymbol{U}^{n}\right\|^{2}+\frac{1}{\bar{\mu}}\left\|\boldsymbol{B}^{n}\right\|^{2}\right)+\left\|\sqrt{\eta^{n-1}} \boldsymbol{D}\left(\boldsymbol{U}^{n}\right)\right\|^{2}+\beta_{1} h^{-\beta_{1}}\left\|\operatorname{div} \boldsymbol{U}^{n}\right\|^{2} \\
& +\beta_{3} h^{\beta_{3}}\left\|\Delta_{h} \boldsymbol{U}^{n}\right\|^{2}+\frac{1}{\bar{\mu}^{2}}\left\|\frac{1}{\sqrt{\xi^{n-1}}} \operatorname{curl} \boldsymbol{B}^{n}\right\|^{2}+\frac{k}{2}\left(\left\|\sqrt{\rho_{+}^{n-1}} d_{t} \boldsymbol{U}^{n}\right\|_{*}^{2}+\beta_{2} h^{\beta_{2}}\left\|\nabla d_{t} \boldsymbol{U}^{n}\right\|^{2}+\left\|d_{t} \boldsymbol{B}^{n}\right\|^{2}\right) \\
0= & \frac{1}{2} d_{t}\left\|\rho^{n}\right\|_{h}^{2}+\frac{k}{2}\left\|d_{t} \rho^{n}\right\|_{h}^{2}+\alpha h^{\alpha}\left\|\nabla \rho^{n}\right\|^{2}
\end{aligned}
$$

Let $V_{h} \cap L_{0}^{2} \subseteq L_{h}, \mathcal{T}_{h}$ be a strongly acute triangulation, $\alpha, \beta_{>} 0$ and $0<\alpha+\frac{\beta_{2}}{2}<\frac{6-d}{6}$. Then $0<\bar{\rho}_{1} \leq \rho^{n} \leq \bar{\varrho}_{2}<\infty$ (discrete maximum principle).

## Sketch of the proof

By asumption the discrete inf-sup-cond. is fullfilled $\Rightarrow \exists R^{n}, P^{n}$.

## Lemma (Existence, energy law, maximum principle, [Baňas and Prohl, 2010])

Let $\alpha, \beta_{1}, \beta_{2}, \beta_{3} \geq 0, \sqrt{\beta_{2}} h^{\beta_{2}} 2\left\|\nabla \boldsymbol{U}^{0}\right\| \leq C$. Then there exists a solution $\left(\boldsymbol{U}^{n}, \boldsymbol{B}^{n}, \rho^{n}, P^{n}, R^{n}\right)$ of the numerical scheme, which holds the discrete energy law

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\begin{aligned}
\int_{\Omega} \rho^{n-1} \boldsymbol{g}^{n} \boldsymbol{U}^{n}= & \frac{1}{2} \partial_{t}\left(\left\|\sqrt{\rho_{+}^{n}} \boldsymbol{U}^{n}\right\|_{*}^{2}+\beta_{2} h^{\beta_{2}}\left\|\nabla \boldsymbol{U}^{n}\right\|^{2}+\frac{1}{\bar{\mu}}\left\|\boldsymbol{B}^{n}\right\|^{2}\right)+\left\|\sqrt{\eta^{n-1}} \boldsymbol{D}\left(\boldsymbol{U}^{n}\right)\right\|^{2}+\beta_{1} h^{-\beta_{1}}\left\|\operatorname{div} \boldsymbol{U}^{n}\right\|^{2} \\
& +\beta_{3} h^{\beta_{3}}\left\|\Delta_{h} \boldsymbol{U}^{n}\right\|^{2}+\frac{1}{\bar{\mu}^{2}}\left\|\frac{1}{\sqrt{\xi^{n-1}}} \operatorname{curl} \boldsymbol{B}^{n}\right\|^{2}+\frac{k}{2}\left(\left\|\sqrt{\rho_{+}^{n-1}} d_{t} \boldsymbol{U}^{n}\right\|_{*}^{2}+\beta_{2} h^{\beta_{2}}\left\|\nabla d_{t} \boldsymbol{U}^{n}\right\|^{2}+\left\|d_{t} \boldsymbol{B}^{n}\right\|^{2}\right) \\
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## Sketch of the proof

Use $\psi:=\frac{1}{\bar{\mu}} \boldsymbol{B}^{n}$ and $\boldsymbol{W}:=\boldsymbol{U}^{n}$ in (Scheme A) and direct calculation.

## Lemma (Existence, energy law, maximum principle, [Bañas and Prohl, 2010])

Let $\alpha, \beta_{1}, \beta_{2}, \beta_{3} \geq 0, \sqrt{\beta_{2}} h^{\beta_{2}} 2\left\|\nabla \boldsymbol{U}^{0}\right\| \leq \boldsymbol{C}$. Then there exists a solution $\left(\boldsymbol{U}^{n}, \boldsymbol{B}^{n}, \rho^{n}, P^{n}, R^{n}\right)$ of the numerical scheme, which holds the discrete energy law

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\begin{aligned}
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$0<\bar{\rho}_{1} \leq \rho^{n} \leq \bar{\varrho}_{2}<\infty$ (discrete maximum principle).

## Sketch of the proof

- Write (Scheme A) as $\mathcal{A}^{n} \boldsymbol{U}^{n}=\mathcal{G}^{n}$ and show that $\mathcal{A}^{n}$ is a M-Matrix.
- By the strongly acute triangulation the off-diagonal entries of the stiffness matrix are negative. Use standard techniques of FEM theory (e.g. inverse estimates) and assumption of constants to proof this.
- Lower bound for $\rho^{n}$ follows from M-Matrix property.
- Upper bound for $\rho^{n}$ : direct calculation.


## Passing to the (weak) limit

Let $\mathcal{U}, \mathcal{B}$, etc. be the (time-) interpolant of $\boldsymbol{U}^{n}, \boldsymbol{B}^{n}$, etc. and all asumptions hold from the last slide. There exists

```
u}\in\mp@subsup{L}{}{\infty}(0,T;\mp@subsup{L}{}{2}\cap{\operatorname{div}\boldsymbol{u}=0\mathrm{ weakly }) ) L'2}(0,T;\mp@subsup{\boldsymbol{W}}{0}{1,2}\cap{\operatorname{div}\boldsymbol{u}=0\mathrm{ a.e. }),
b}\in\mp@subsup{L}{}{\infty}(0,T;\mp@subsup{\boldsymbol{L}}{}{2}\cap{\operatorname{div}\boldsymbol{u}=0\mathrm{ weakly }) }\cap\mp@subsup{L}{}{2}(0,T;\boldsymbol{X})\mathrm{ and }\rho\in\mp@subsup{L}{}{\infty}(0,T;\mp@subsup{L}{}{\infty})\mathrm{ s.t.
```

$$
\begin{aligned}
\mathcal{U} \rightarrow^{*} \boldsymbol{u} & \text { in } L^{\infty}\left(0, T ; \boldsymbol{L}^{2}\right), \\
\mathcal{U} \rightarrow \boldsymbol{u} & \text { in } L^{2}\left(0, T ; \boldsymbol{W}^{1,2}\right), \\
\operatorname{div} \mathcal{U} \rightarrow 0 & \text { in } L^{2}\left(0, T ; L^{2}\right)\left(\beta_{1}>0\right), \\
\mathcal{B} \rightarrow^{*} \boldsymbol{b} & \text { in } L^{\infty}\left(0, T ; \boldsymbol{L}^{2}\right), \\
\mathcal{B} \rightarrow^{2} \boldsymbol{b} & \text { in } L^{2}(0, T ; \boldsymbol{H}(\text { curl })), \\
\sigma \rightarrow^{*} \rho & \text { in } L^{\infty}\left(0, T ; L^{\infty}\right) .
\end{aligned}
$$

$\boldsymbol{X}:=\boldsymbol{H}($ curl $) \cap \boldsymbol{H}_{0}($ div $) \cap\{\operatorname{div} \boldsymbol{b}=0$ a.e. $\}$

## Passing to the (weak) limit

Let $\mathcal{U}, \mathcal{B}$, etc. be the (time-) interpolant of $\boldsymbol{U}^{n}, \boldsymbol{B}^{n}$, etc. and all asumptions hold from the last slide. There exists
$\boldsymbol{u} \in L^{\infty}\left(0, T ; L^{2} \cap\{\operatorname{div} \boldsymbol{u}=0\right.$ weakly $\left.\}\right) \cap L^{2}\left(0, T ; \boldsymbol{W}_{0}^{1,2} \cap\{\operatorname{div} \boldsymbol{u}=0\right.$ a.e. $\left.\}\right)$,
$\boldsymbol{b} \in L^{\infty}\left(0, T ; \boldsymbol{L}^{2} \cap\{\operatorname{div} \boldsymbol{u}=0\right.$ weakly $\left.\}\right) \cap L^{2}(0, T ; \boldsymbol{X})$ and $\rho \in L^{\infty}\left(0, T ; L^{\infty}\right)$ s.t.

$$
\begin{aligned}
\mathcal{U} \rightarrow^{*} \boldsymbol{u} & \text { in } L^{\infty}\left(0, T ; \boldsymbol{L}^{2}\right), \\
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\operatorname{div} \mathcal{U} \rightarrow 0 & \text { in } L^{2}\left(0, T ; L^{2}\right)\left(\beta_{1}>0\right), \\
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$$

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## Sketch of the proof

- Boundness of follows direct from the energy law.
- $\operatorname{div} \boldsymbol{u}=0$ a.e. follows directly.
- As $\operatorname{div} \mathcal{B}=0$, we have $\operatorname{div} \boldsymbol{b}=0$ by a result of Kikuchi, cf. [Hiptmair, 2002, Theorem 4.9].

In order to gain convergence results, we need strong $L^{2}$-convergence of $\rho$ !

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## Lemma

Under certain assumptions for the constants and initial data, $\mathcal{U}^{+} \rightarrow \boldsymbol{u}$ in $L^{2}\left(0, T ; L^{2}\right), \sigma^{+} \rightharpoonup^{*} \rho$ in $L^{\infty}\left(0, T ; L^{2}\right)$, we have

- $\rho$ is unique weak solution of $\rho_{t}+\operatorname{div}(\rho \boldsymbol{u})=0$,
- $\sigma \rightarrow \rho$ in $L^{2}\left(0, T ; L^{2}\right)$.

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- $\rho$ is unique weak solution of $\rho_{t}+\operatorname{div}(\rho \boldsymbol{u})=0$,
- $\sigma \rightarrow \rho$ in $L^{2}\left(0, T ; L^{2}\right)$.


## Sketch of the proof

- Weak convergence is assured by last lemma. Show that $\left\|\sigma^{+}\right\| \rightarrow\|\rho\|$.
- Use (variant (cf. [Walkington, 2004]) of) DiPerna-Lions compactness argument ([DiPerna and Lions, 1989]) and standard arguments (weak continuity of the norm, assumptions, Fatou, energy law, etc.) to conclude

We need strong convergence of $\boldsymbol{u}$ and $\boldsymbol{b}$ in $L^{2}\left(0, T ; \boldsymbol{L}^{2}\right)$ !

We need strong convergence of $\boldsymbol{u}$ and $\boldsymbol{b}$ in $L^{2}\left(0, T ; \boldsymbol{L}^{2}\right)$ !

## Lemma (Compactness result, [Lions and Magenes, 1972])

If there exists $C>0, \alpha>0$, s.t. for all $0<\delta \leq T$

$$
\int_{\delta}^{T}\left\|v_{h}(t)-v_{h}(t-\delta)\right\|_{L^{2}} \leq C \delta^{\alpha},
$$

then, the seq. $\left(v_{h}\right)$ is compact in $L^{2}\left(0, T ; L^{2}\right)$.

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## Lemma ([Baňas and Prohl, 2010])

Let $(., .)_{*}=(., .)_{h}, V_{h} \cap L^{2} \subseteq L_{h}$ and $0<\alpha, \beta_{1}, \beta_{2}, \beta_{3}$ s.t. $0<\alpha+\frac{\beta_{2}}{2}<\frac{6-d}{6}$ and $\alpha \geq \frac{\beta_{3}}{4}$ (for $d=3$ ). Then there exists $C>0, \alpha>0$ s.t.

$$
\int_{\delta}^{T}\|\mathcal{U}(t, .)-\mathcal{U}(t-\delta, .)\|_{L^{2}}^{2}+\|\mathcal{B}(t, .)-\mathcal{B}(t-\delta, .)\|_{L^{2}}^{2} d t \leq C \delta^{k} \quad \forall \delta \in[0, T] .
$$

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Let $(., .)_{*}=(., .)_{h}, V_{h} \cap L^{2} \subseteq L_{h}$ and $0<\alpha, \beta_{1}, \beta_{2}, \beta_{3}$ s.t. $0<\alpha+\frac{\beta_{2}}{2}<\frac{6-d}{6}$ and $\alpha \geq \frac{\beta_{3}}{4}$ (for $d=3$ ). Then there exists $C>0, \alpha>0$ s.t.

$$
\int_{\delta}^{T}\|\mathcal{U}(t, .)-\mathcal{U}(t-\delta, .)\|_{L^{2}}^{2}+\|\mathcal{B}(t, .)-\mathcal{B}(t-\delta, .)\|_{L^{2}}^{2} d t \leq C \delta^{k} \quad \forall \delta \in[0, T] .
$$

## Sketch of the proof

- much direct, technical calculation. Frequently use of energy law, standard estimates of FEM theory.
- For the $\mathcal{B}$ terms use of property of Hodge map, cf. [Hiptmair, 2002].

We need strong convergence of $\boldsymbol{u}$ and $\boldsymbol{b}$ in $L^{2}\left(0, T ; \boldsymbol{L}^{2}\right)$ !

## Lemma (Compactness result, [Lions and Magenes, 1972])

If there exists $C>0, \alpha>0$, s.t. for all $0<\delta \leq T$

$$
\int_{\delta}^{T}\left\|v_{h}(t)-v_{h}(t-\delta)\right\|_{L^{2}} \leq C \delta^{\alpha}
$$

then, the seq. $\left(v_{h}\right)$ is compact in $L^{2}\left(0, T ; L^{2}\right)$.

## Lemma ([Baňas and Prohl, 2010])

Let $(., .)_{*}=(., .)_{h}, V_{h} \cap L^{2} \subseteq L_{h}$ and $0<\alpha, \beta_{1}, \beta_{2}, \beta_{3}$ s.t. $0<\alpha+\frac{\beta_{2}}{2}<\frac{6-d}{6}$ and $\alpha \geq \frac{\beta_{3}}{4}$ (for $d=3$ ). Then there exists $C>0, \alpha>0$ s.t.

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## Corollary

The weak limits $\mathbf{u}, \boldsymbol{b}, \rho$ and $p, R$ solve the weak two-fluid MHD equation.

## Discontinuous Galerkin Approach

In [Liu and Walkington, 2007] they propose a discontinuous Galerkin scheme for the denstity depend Navier-Stokes eq. with p.w. constant FE space w.r.t. $\rho$ and aver. divergence zero w.r.t. $u$.

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## Continuous Galerkin

## Advantage

- gen: direct use of appropriate testfunctions $\Rightarrow$ energy law.


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## Advantage

- gen: Possibility of higher polynomial degrees.
- gen: Flexibility in grid design.


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- gen: High grow of dof for computition, bad condition number of stiffness matrix.


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- gen: prob. high effort near the boundary.


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## Continuous Galerkin

## Advantage

- gen: direct use of appropriate testfunctions $\Rightarrow$ energy law.
- here: M-matrix property.
- here: here: Taylor-Hood, MINI elements are admitted.


## Disadvantage

- gen: prob. high effort near the boundary.
- here: stabilazation terms required for convergence.


## Discontinuous Galerkin

## Advantage

- gen: Possibility of higher polynomial degrees.
- gen: Flexibility in grid design.
- here: Monotonicity of the iterates.


## Disadvantage

- gen: High grow of dof for computition, bad condition number of stiffness matrix.
- here: control of jump terms difficult.


## Outline

## (1) Introduction, prelimitaries <br> (2) Discretization, Convergence

(3) Some words on implementation
(4) Summary

- Solve (Scheme A) by a fixed point scheme which decouples $\boldsymbol{u}$ and $\boldsymbol{b}$ (right hand side of the eq. old iterate). This scheme converges to weak solution of the two-fluid MHD eq.
- Implementation of NSE with Taylor-Hood elements.


Figure 2. Density at $x=0.5$ at times $t=0,4,14,20$.


## Outline

(9) Introduction, prelimitaries
(2) Discretization, Convergence
(5) Some words on implementation
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## Summary

- Discretization based on FEM (via continuous Galerkin method!) and Euler scheme for time. Calculation simplifies with splitting scheme.
- No convergence rates!
- Mostly advatages in contrast to discontinuous Galerkin methods (M-matrix property, no jump terms, etc.)


## Thanks

## Thank you for your attention!

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